Contents lists available at ScienceDirect

# Journal of Nuclear Materials

journal homepage: www.elsevier.com/locate/jnucmat

# First-principles study of structural, elastic and electronic properties of thorium dicarbide (ThC<sub>2</sub>) polymorphs

# I.R. Shein, A.L. Ivanovskii \*

Institute of Solid State Chemistry, Ural Branch of the Russian Academy of Sciences, Pervomaiskaya Str. 91, Ekaterinburg 620041, Russia

#### ARTICLE INFO

Article history: Received 19 February 2009 Accepted 2 June 2009

#### ABSTRACT

The comparative study of the structural, elastic, cohesive and electronic properties of three polymorphs ( $\alpha$ -monoclinic,  $\beta$ -tetragonal and  $\gamma$ -cubic) of thorium dicarbide ThC<sub>2</sub> is performed within the density-functional theory. The optimized atomic coordinates, lattice parameters, theoretical density ( $\rho$ ), bulk moduli (B), compressibility ( $\beta$ ), as well as electronic densities of states, electronic heat capacity ( $\gamma$ ) and molar Pauli paramagnetic susceptibility ( $\chi$ ) for all ThC<sub>2</sub> polymorphs are obtained and analyzed in comparison with available experimental data. The peculiarities of inter-atomic bonding for thorium dicarbide are discussed. Besides, we have evaluated the formation energies ( $E_f$ ) of ThC<sub>2</sub> polymorphs for different possible preparation routes (namely for the reactions with the participation of simple substances (metallic Th and graphite) or thorium monocarbide ThC and graphite). The results show that the synthesis of the ThC<sub>2</sub> polymorphs from simple substances is more favorable – in comparison with the reactions with participation of Th monocarbide.

© 2009 Elsevier B.V. All rights reserved.

# 1. Introduction

The binary compounds formed by thorium with light *sp* elements possess high density, good thermal conductivity and some other interesting physical properties which makes them attractive from a fundamental point of view as well as for a variety of technical applications for example as alternative fertile materials to be used in nuclear breeder systems *etc*, see [1-4]. The thorium carbides belong to this important group of materials.

Thorium reacts with carbon to form the two main phases. One of them is thorium monocarbide (ThC), which crystallizes in the cubic *B*1-like structure, where carbon atoms occupy the octahedral interstitial sites in the face-centered cubic thorium sublattice. This phase is stable over a wide range of C/Th compositions (ThC<sub>x</sub> where  $1.0 \ge x > 0.55$  [3]), *i.e.* can possess the high carbon deficiency.

The second phase is carbon-rich thorium dicarbide ThC<sub>2</sub>, with relatively narrow range of homogeneity. Thorium dicarbide exists as stoichiometric (C/Th = 2) phase in three polymorphic modifications: low-temperature  $\alpha$ -ThC<sub>2</sub> polymorph (up to 1700 K) and two high-temperature polymorphs:  $\beta$ -ThC<sub>2</sub> (in range 1700 K < *T* < 1754 K) and  $\gamma$ -ThC<sub>2</sub> (at *T* > 1754 K), see [3,5,6].

The physical properties and chemical behavior exhibited by the thorium carbides are closely related to their electronic structure. Today, a set of the theoretical studies have been performed to understand the electronic structure and the origin of properties of thorium monocarbide ThC [3,7–13]. As distinct from the cubic

Th monocarbide, the data on the fundamental electronic properties for ThC<sub>2</sub> polymorphs are very limited. To our knowledge, up to now only the structural and electronic parameters for low-temperature modification  $\alpha$ -ThC<sub>2</sub> were studied theoretically using the densityfunctional theory [11].

In this work. we present the results of *ab* initio study of the structural, elastic, and electronic properties of all three  $\alpha$ ,  $\beta$  and  $\gamma$  ThC<sub>2</sub> polymorphs. As a result, the equilibrium geometries, theoretical density ( $\rho$ ), bulk moduli (*B*) and compressibility ( $\beta$ ), electronic densities of states (DOS), the Sommerfeld constant ( $\gamma$ ) and the Pauli paramagnetic susceptibility ( $\chi$ ) for all mentioned polymorphs are obtained for the first time and analyzed in comparison with available experimental data. Additionally, the theoretical estimations of the formation energies for ThC<sub>2</sub> polymorphs which can provide the important insights of the fundamental aspects of phase equilibrium in the Th–C system are performed. The peculiarities of the chemical bonding in thorium dicarbide are also discussed.

# 2. Models and method of calculation

The common structural feature of all  $ThC_2$  polymorphs is that these materials include in their crystal frameworks the discrete  $C_2$  structural units – so-called  $C_2$  dumbbells, see Fig. 1.

The low-temperature  $\alpha$  modification of ThC<sub>2</sub> has a monoclinic structure (space group C2/*c*) with four formula units in elementary cell (*Z* = 4); the cell structure and the orientation of the C<sub>2</sub> dumbbells are illustrated in Fig. 1. The intermediate  $\beta$  modification of ThC<sub>2</sub> has a body-centered tetragonal structure (space group





<sup>\*</sup> Corresponding author. Tel.: +7 343 3755331; fax: +7 343 3744495. *E-mail address:* ivanovskii@ihim.uran.ru (A.L. Ivanovskii).

<sup>0022-3115/\$ -</sup> see front matter @ 2009 Elsevier B.V. All rights reserved. doi:10.1016/j.jnucmat.2009.06.004



Fig. 1. Atomic structures of ThC<sub>2</sub> polymorphs: (1) monoclinic  $\alpha$ -ThC<sub>2</sub>; (2) tetragonal  $\beta$ -ThC<sub>2</sub> and (3) cubic  $\gamma$ -ThC<sub>2</sub>.

I4/*mmm*, *Z* = 2) where the C<sub>2</sub> dumbbells are oriented along the *c* axis. Finally, the high-temperature  $\gamma$  modification of ThC<sub>2</sub> crystallizes in a face-centered cubic structure (KCN-structural type, space group Fm3m, *Z* = 4), see Fig. 1.

The band-structure calculations of the all mentioned ThC<sub>2</sub> polymorphs were done by means of the density-functional theory (DFT) and the Vienna *ab initio* simulation package (VASP) [14,15] in projector augmented waves (PAW) formalism [16,17] and spin–orbit interactions. Exchange and correlation were described by a nonlocal correction for LDA [18] in the form of the generalized gradient approximation (GGA) [19]. The kinetic energy cutoff of 600 eV and *k*-mesh of  $8 \times 8 \times 8$  were used. The geometry optimization was performed with the force cutoff of 1 meV/Å.

## 3. Results and discussion

# 3.1. Structural and elastic properties

As the first step, the total energy ( $E_{tot}$ ) versus cell volume calculations were carried out to determine the equilibrium structural parameters for considered ThC<sub>2</sub> polymorphs and their relative stability. These calculations were performed with full-lattice optimization including atomic coordinates. The self-consistent calculations were considered to be converged when the difference in the total energy of the crystal did not exceed 0.01 meV as calculated at consecutive steps.

Fig. 2 shows the variation for  $E_{tot}$  from cell volume V; the calculated equilibrium lattice constants and atomic positions are given



**Fig. 2.** Energy minimization with respect to cell volumes for ThC<sub>2</sub> polymorphs: monoclinic  $\alpha$ -ThC<sub>2</sub>, tetragonal  $\beta$ -ThC<sub>2</sub> cubic  $\gamma$ -ThC<sub>2</sub>. The values of the total energies are given relative to the  $E_{tot}$  of the most stable  $\alpha$ -ThC<sub>2</sub> polymorph.

in Table 1. The results obtained allow us to make the following conclusions.

(1) In agreement with experiments [3], our total energy calculations show that the most stable polymorph is monoclinic  $\alpha$ -ThC<sub>2</sub>, whereas  $\beta$ -ThC<sub>2</sub> is intermediate and cubic  $\gamma$ -ThC<sub>2</sub> is the most unstable polymorph and presumably can be stabilizing at high temperatures by the entropy factor. In order to estimate the relative stability of ThC<sub>2</sub> polymorphs, their Gibbs free energies ( $G = E^{\circ} + P \cdot V - T \cdot S$ ) should be obtained. Since all our calculations are performed at zero temperature and zero pressure conditions, the G becomes equal to the  $E^{o}$ and, for example, for comparative estimation of stability of  $\alpha$ and  $\beta$  polymorphs the value of  $\Delta E_{tot} = \{E_{tot}(\alpha - ThC_2) - C_2\}$  $E_{tot}(\beta-ThC_2)$  may be used. In this way, we have obtained  $\Delta E_{tot} = 0.063 \text{ eV/form.unit}$  (6.08 kJ/mol) for  $\alpha$ -ThC<sub>2</sub>  $\leftrightarrow \beta$ -ThC<sub>2</sub> transition and  $\Delta E_{tot} = 1.707 \text{ eV/form.unit}$  (164.72 kJ/ mol) for  $\beta$ -ThC<sub>2</sub>  $\leftrightarrow \gamma$ -ThC<sub>2</sub> transition. In the experiments, the enthalpy for  $\alpha$ -ThC<sub>2</sub>  $\leftrightarrow \beta$ -ThC<sub>2</sub> and  $\beta$ -ThC<sub>2</sub>  $\leftrightarrow \gamma$ -ThC<sub>2</sub> transitions are estimated at about 2.1 and 10.5 kJ/mol, respectively [3]. Naturally, the thermodynamic factors such as temperature and pressure should be taken into account at such comparisons.

#### Table 1

Calculated total energies ( $E_{tot}$ , eV per ThC<sub>2</sub> molecule), atomic positions, lattice constants (a, b and c, in Å), d(C–C) distances (for C<sub>2</sub> dumbbells, in Å), and theoretical density ( $\rho$ , in g/cm<sup>3</sup>) for ThC<sub>2</sub> polymorphs in comparison with available experimental data.

Polymorph/ parameter	α-ThC <sub>2</sub>	β-ThC <sub>2</sub>	$\gamma$ -ThC <sub>2</sub>
Etot <sup>*</sup>	0.000	0.063	1.770
а	6.6283 (6.692 [5] 6.684 [6])**	3.8069 (4.221 [5]	5.8097
		4.235 [20])	(5.806 [5]
			5.807 [20])
b	4.2069 (4.223 [5] 4.220 [6])	-	-
с	6.6726 (6.744 [5] 6.735 [6])	6.3826 (5.395 [5]	-
		5.408 [20])	
Wyckoff	$4(e) 0; 0.1990; \frac{1}{4} (0; 0.2074;$	2(a) 0; 0; 0	4(a) 0; 0; 0
positions for Th	<sup>1</sup> ⁄ <sub>4</sub> ) [5] (0; 0.2022; <sup>1</sup> ⁄ <sub>4</sub> ) [6]		
Wyckoff	8(f) 0.3006; 0.1356; 0.0523	4(e) 0; 0; 0.3957	4(b) 0; 0;
positions	(0.2992; 0.1326; 0.0540) [5];	(0; 0; ½ [5,18] -	0.3375 (½;
for	(0.2986; 0.1305; 0.0513) [6]	for dimer C <sub>2</sub> )	1/2; 1/2 [5,20] -
carbon			for dimer C <sub>2</sub> )
d(C-C)	1.328	1.331	1.326
ρ	9.03 (9.14 [3])****	9.15	8.64 (9.00 [3]

<sup>\*</sup> With respect to the  $E_{tot}$  of  $\alpha$ -ThC<sub>2</sub>.

\*\* Available experimental data are given in parentheses.

For composition  $\alpha$ -ThC<sub>1.94</sub> [3].

- (2) The calculated structural parameters for all ThC<sub>2</sub> polymorphs are summarized in Table 1. As can be seen, for  $\alpha$ and  $\gamma$  polymorphs our results are very similar to experiments. On the other hand, for tetragonal  $\beta$ -ThC<sub>2</sub> our calculations have overestimated the parameter c ( $c^{\text{theor}} = 6.3826 \text{ Å}$ *versus*  $c^{exp} = 5.395 - 5.408 \text{ Å} [5,20]$ ) whereas, the parameter *a* is underestimated ( $a^{\text{theor}} = 3.8069 \text{ Å}$  versus  $a^{\text{exp}} = 4.221 - 4.221$ 4.235 Å [5,20]). This result presumably may be explained taking into account the partial thermal disorientation for C<sub>2</sub> dumbbells (for example, the orientation of a part of the  $C_2$  dumbbells in the plane perpendicular to the *c* axis, see Ref. [3]) – in comparison with ideal ordering of C<sub>2</sub> dumbbells along *c* axis, as it was postulated in our model. Besides it is necessary to take into account the non-stoichiometry effect: the tetragonal  $\beta$ -ThC<sub>2</sub> has a homogeneity range between C/Th = 1.66 and 1.94 [3].
- (3) The bulk moduli (*B*) and its pressure derivative (B'), derived from the fit to Murnaghan's equation of state, as well as compressibility ( $\beta = 1/B$ ), are given in Table 2. From these data, we see that the maximal and minimal values of the bulk moduli adopt  $\beta$ -ThC<sub>2</sub> and  $\gamma$ -ThC<sub>2</sub>; whereas the compressibility of ThC<sub>2</sub> polymorphs changes in the inverse sequence. A simple explanation of these data follows from the known correlation between the bulk moduli *B* and cell volumes  $V (B \sim V^{-k} [21])$ , see Table 2. Thus, among all ThC<sub>2</sub> polymorphs,  $\gamma$ -ThC<sub>2</sub> adopts the minimal *B* and maximal compressibility; i.e. in comparison with other ThC2 polymorphs,  $\gamma$ -ThC<sub>2</sub> is the most soft material. Finally, the comparison of our data (Table 2) with the same for cubic ThC monocarbide [3,8,10] demonstrate that  $B(ThC_2) > B(ThC)$ . Thus, based on the well-known correlations between the bulk moduli and mechanical strength [22], we conclude that with growth of carbon content (i.e. as going from ThC to ThC<sub>2</sub>) the hardening of a material occur. This conclusion agree with Vickers hardness measurements for ThC<sub>z</sub> samples with the variable Th/C content (1 < z < 2), see [3].

# 3.2. Energies of formation

To provide the further insight of the fundamental aspects of phase equilibrium in the Th–C system, we have performed the estimations of the formation energies for ThC<sub>2</sub> polymorphs. For this purpose, the total energies of ThC<sub>2</sub> and its constituents in formal reactions ThC<sub>2</sub> = Th + 2C<sup>g</sup> or ThC<sub>2</sub> = ThC + C<sup>g</sup> are calculated, and the formation energies ( $E_f$ ) of ThC<sub>2</sub> polymorphs are defined as:

$$E_{\rm f}({\rm Th}C_2) = E_{\rm tot}({\rm Th}C_2) - \{E_{\rm tot}({\rm Th}) + 2E_{\rm tot}({\rm C}^{\rm g})\} \tag{1}$$

$$E_{\rm f}({\rm Th}C_2) = E_{\rm tot}({\rm Th}C_2) - \{E_{\rm tot}({\rm Th}C) + E_{\rm tot}({\rm C}^{\rm g})\} \tag{2}$$

where  $E_{tot}$  are the total energies of ThC<sub>2</sub> polymorphs, cubic ThC, metallic thorium ( $\alpha$ -Th), and graphite (C<sup>g</sup>) at their optimized geometries as calculated in the same VASP approach. In this way, negative values of  $E_f$  point out the formation of stable dicarbide polymorphs, but their formation will be unlikely at the ordinary conditions if  $E_f > 0$ . Certainly, this prediction is based only on the

**Table 2** Calculated bulk moduli (*B*, in GPa), their energy derivatives (*B'*), compressibility ( $\beta$ , in 1/GPa), and cell volumes (*V*, in Å<sup>3</sup>) for ThC<sub>2</sub> polymorphs.

Polymorph/parameter	$\alpha$ -ThC <sub>2</sub>	β-ThC <sub>2</sub>	$\gamma$ -ThC <sub>2</sub>
В	129.1	149.2	121.6
Β'	3.84	4.13	3.71
β	0.00775	0.00670	0.00822
V	46.92	46.28	49.02

thermodynamics and does not take into account the kinetics of reactions.

Formation energies calculated for the different reactions are shown in Table 3 and allow us to make the following conclusions.

- (1) For  $\alpha$  and  $\beta$  polymorphs, we obtained, that their formation from the simple reagents ( $\alpha$ -Th and C<sup>g</sup>) as well as from ThC and carbon is favorable – against  $\gamma$  polymorph, for which the formation energy was obtained as positive.
- (2) For both reactions (1) and (2), we have obtained, that the negative values of  $E_{\rm f}$  for  $\alpha$  polymorph are higher than for  $\beta$  polymorph; this means that the stability of  $\alpha$ -ThC<sub>2</sub> is higher in accordance with the thermodynamic data [3].
- (3) Synthesis of  $ThC_2$  with the participation of ThC as a reagent is the most energetically unfavorable, and this fact can be explained by the high stability of thorium monocarbide [3,10].

# 3.3. Electronic properties

In order to gain a deeper insight of the comparative peculiarities of the electronic properties of ThC<sub>2</sub>, the calculations of their electronic structures have been performed. Total densities of states (DOS) for all ThC<sub>2</sub> polymorphs, depicted in Fig. 3, contain three main bands (A–C). The most significant differences for ThC<sub>2</sub> polymorphs are the width of the occupied valence bands (~7.6 eV for  $\beta$ -, ~8.6 eV for  $\gamma$ - and ~7.3 eV for  $\alpha$ -polymorph) and the presence of the gap (~0.4 eV) between A and B bands for monoclinic  $\alpha$ -ThC<sub>2</sub> *versus* the continuous character of spectra for the others polymorphs.

The total density of states at the Fermi level  $N(E_F)$  adopts the non-zero value (*i.e.* all polymorphs have a metallic conductivity) and the minimal value of  $N(E_F)$  is for the most stable monoclinic  $\alpha$ -polymorph, Table 4.

The above calculations allow us to estimate for ThC<sub>2</sub> polymorphs the specific heat coefficient ( $\gamma$ ) and the Pauli paramagnetic susceptibility ( $\chi$ ), assuming the free electron model, as:  $\gamma = (\pi^2/3)N(E_F)k_B^2$ , and  $\chi = \mu_B^2N(E_F)$ ; the results show the tendency

# Table 3

Formation energies ( $E_{\rm f}$  eV/Th atom) for ThC<sub>2</sub> polymorphs for the different reactions.

Polymorph/reaction	$\alpha$ -ThC <sub>2</sub>	β-ThC <sub>2</sub>	γ-ThC <sub>2</sub>
ThC <sub>2</sub> = Th + 2C <sup>g</sup>	-0.92	-0.86	0.27
ThC <sub>2</sub> = ThC + C <sup>g</sup>	-0.22	-0.16	0.97



Fig. 3. Total densities of states for  $ThC_2$  polymorphs: (1) tetragonal  $\beta$ -ThC<sub>2</sub>; (2) cubic  $\gamma$ -ThC<sub>2</sub> and (3) monoclinic  $\alpha$ -ThC<sub>2</sub>.

of growth of  $\gamma$  and  $\chi$  as going from  $\alpha\text{-ThC}_2$  to  $\beta,\gamma\text{-ThC}_2$ . Let us note also that the calculated specific heat coefficient  $(\gamma)$  for stable  $\alpha\text{-ThC}_2$  (2.614 mJ mol $^{-1}$  K $^{-2}$ ) is higher than the same for cubic monocarbide ThC  $(\gamma^{ThC} \sim 1.87 \text{ mJ mol}^{-1} \text{ K}^{-2})$  [8]; and this tendency agree with the available experiments:  $\gamma^{ThC}_2 = 3.13 \text{ mJ} \text{ mol}^{-1} \text{ K}^{-2} > \gamma^{ThC} = 2.12 - 2.92 \text{ mJ mol}^{-1} \text{ K}^{-2}$  [3].

# 3.4. Chemical bonding

Let us discuss the peculiarities of chemical bonding in thorium dicarbides. Further for the analysis of the hybridization effects (*i.e. the covalent bonding* character) the atomic-resolved *l*-projected DOS, and the charge density maps are employed. For the description of *ionic bonding* for ThC<sub>2</sub> we have made a Bader [23] analysis. In this approach each atom of a crystal is surrounded by an effective surface that runs through minima of the charge density, and the total charge of an atom (so-called Bader charge,  $Q^B$ ) is determined by integration within this region.

Formally, one could consider  $\text{Th}C_2$  crystal as  $\text{Th}^{4+}(C_2)^{4-}$ , where thorium is as four-electron donor. However, this simplified 'pure ionic' picture ignores the presence of covalent thorium–carbon bonding.

As can be seen from the Fig. 4, where on the example of  $\beta$ -ThC<sub>2</sub> polymorph the site-projected DOSs are depicted, there are significant hybridization effects of C 2*p* states with the Th (6*d* + 5*f*) states

#### Table 4

Total densities of states at the Fermi level (N( $E_F$ ), in states/eV-atom Th), electronic heat capacity  $\gamma$  (in mJ K<sup>-2</sup> mol<sup>-1</sup>) and molar Pauli paramagnetic susceptibility  $\chi$  (in 10<sup>-4</sup> emu/mol) for ThC<sub>2</sub> polymorphs.

Polymorph/parameter	$\alpha$ -ThC <sub>2</sub>	β-ThC <sub>2</sub>	$\gamma$ -ThC <sub>2</sub>
$N(E_F)$	1.109	1.426	1.410
γ	2.614	3.361	3.324
χ	0.357	0.459	0.454



Fig. 4. Total (upper panel) and partial atomic-resolved densities of states for tetragonal  $\beta$ -ThC<sub>2</sub> polymorph.



**Fig. 5.** Valence charge density map in (1 1 0) plane of  $\beta$ -ThC<sub>2</sub> polymorph.

(in the interval from -7.6 up to -2.2 eV), which indicate the presence of the strong covalent Th–C bonding. Let us note also, that the Th–C bonds are formed not only along of the *c* axis (as a quasi-one-dimensional  $\sigma$ -like bonds ..Th–C–C–Th–C–C–Th..), but also 'lateral'  $\pi$ -like overlapping C<sub>2</sub>–Th takes place, see Fig. 5. Here, it can be seen also that 'direct' Th–Th bonds are very small.

Finally, our estimations of Bader charges show that the amounts of electrons redistributed between the Th atoms and  $C_2$  dumbbells are far from those predicted in the idealized ionic model (4*e*); namely, the charge transfer  $\Delta Q(\text{Th} \rightarrow C_2)$  is about 1.95*e*.

# 4. Conclusions

In summary, the first-principles VASP-PAW with spin–orbit calculations have been performed for systematic study of structural, elastic, cohesive and electronic properties of all three polymorphs ( $\alpha$ -monoclinic,  $\beta$ -tetragonal and  $\gamma$ -cubic) of thorium dicarbide ThC<sub>2</sub>. The main conclusions are the following.

According to our calculations, among all ThC<sub>2</sub> polymorphs, high-temperature  $\gamma$ -ThC<sub>2</sub> with maximal cell volume and minimal density adopts the minimal *B* and maximal compressibility; *i.e.* in comparison with other ThC<sub>2</sub> polymorphs,  $\gamma$ -ThC<sub>2</sub> is the most soft material. We conclude also that with growth of carbon content (*i.e.* as going from ThC to  $\alpha$ -ThC<sub>2</sub>) the hardening of a material occur.

The estimations of the formation energies of  $ThC_2$  polymorphs for different possible preparation routes (namely for the reactions with the participation of simple substances (metallic Th and graphite) or thorium monocarbide ThC and graphite) show that the synthesis of the  $ThC_2$  polymorphs from simple substances is more favorable – in comparison with the reactions with participation of Th monocarbide.

The present results indicate that the electronic spectra of all  $ThC_2$  polymorphs are similar, with non-zero value of the density

of states at the Fermi level, *i.e.* all these materials have a metallic conductivity. The minimal value of  $N(E_F)$  is for the most stable monoclinic  $\alpha$  polymorph. Our analysis reveals that bonding in ThC<sub>2</sub> polymorphs is of a mixed covalent-ionic-metallic character. Namely, the covalent bonding is formed due to hybridization effects of C-C states (for  $C_2$  dumbbells) and  $C_2$ -Th states. In addition, between the thorium atoms and  $C_2$  dumbbells, ionic bonds emerge owing to the charge transfer  $Th \rightarrow C_2$ . Note that the amounts of electrons redistributed between the Th atoms and C<sub>2</sub> dumbbells (at about 1.95e) are much less than assumed in the idealized ionic model (4e). The metallic-like Th-Th bonds are formed by near-Fermi delocalized Th *d*,*f* states.

Finally, our analysis was performed for stoichiometric ThC<sub>2</sub> polymorphs with idealized crystal structures. As an important issue for future studies, the great interest is the presence of the lattice vacancies (*i.e.* the effect of the carbon non-stoichiometry) as well as the partial disorientation of the C<sub>2</sub> dumbbells and their comparative influence on the properties of ThC<sub>2</sub> polymorphs.

## References

- [1] A.J. Freeman, G.H. Lander (Eds.), Handbook of the Physics and Chemistry of the Actinides, North-Holland, Amsterdam, 1985.
- [2] R. Benz, A. Naoumidis, Thorium, Compounds with Nitrogen, Gmelin Handbook of Inorganic Chemistry, eighth Ed., Thorium supplement, vol. C3, Springer, Berlin, 1987.

- [3] H. Kleykamp. Thorium Carbides. Gmelin Handbook of Inorganic and Organometallic Chemistry, eighth Ed., Thorium supplement, vol. C6, Springer. Berlin. 1992
- A.H.M. Evensen, R. Catherall, P. Drumm, P. Van Duppen, O.C. Jonsson, E. Kugler, [4] J. Lettry, O. Tengblad, V. Tikhonov, H.L. Ravn, Nucl. Instrum. Meth. Phys. Res. B 126 (1997) 160
- [5] A.L. Bowman, N.H. Krikorian, G.P. Arnold, T.C. Wallace, N.G. Nereson, Acta Crystallogr. B 24 (1968) 1121.
- [6] D.W. Jones, I.J. McColm, R. Steadman, J. Yerkess, J. Solid State Chem. 68 (1987) 219
- [7] T. Das, S. Deb, A. Mookerjee, Physica B 367 (2005) 6.
- I.R. Shein, K.I. Shein, A.L. Ivanovskii, J. Nucl. Mater. 353 (2006) 19. [8]
- [9] I.R. Shein, K.I. Shein, G.P. Shveikin, A.L. Ivanovskii, Doklady Phys. Chem. 407 (2006) 106.
- [10] I.R. Shein, K.I. Shein, A.L. Ivanovskii, Tech. Phys. Lett. 33 (2007) 128.
- [11] I.R. Shein, K.I. Shein, N.I. Medvedeva, A.L. Ivanovskii, Phys. Stat. Sol. B 244 (2007) 3198.
- [12] I.R. Shein, A.L. Ivanovskii, J. Struct. Chem. 49 (2008) 348.
- [13] I.S. Lim, G.E. Scuseria, Chem. Phys. Lett. 460 (2007) 137.
- [14] G. Kresse, J. Hafner, Phys. Rev. B 47 (1993) 558 [15] G. Kresse, J. Furthmuller, Phys. Rev. B 54 (1996) 11169.
- [16] P.E. Blőhl, Phys. Rev. B 50 (1994) 17953.
- [17] G. Kresse, D. Joubert, Phys. Rev. B 59 (1999) 1758. [18] J.P. Perdew, A. Zunger, Phys. Rev. B 23 (1981) 5048.
- [19] J.P. Perdew, Y. Wang, Phys. Rev. B 45 (1992) 13244.
- [20] P.K. Gantzel, S. Langer, N.L. Baldwin, F.L. Kester, Adv. X-Ray Anal. 8 (1965) 78.
- [21] A.J. Majewski, P. Vogl, Phys. Rev. B 35 (1987) 9666.
- [22] J. Haines, J.M. Leger, G. Bocquillon, Annu. Rev. Mater. Res. 31 (2001) 1.
- [23] R.F.W. Bader, Atoms in Molecules: A Quantum Theory, International Series of Monographs on Chemistry, Clarendon Press, Oxford, 1990.